

# Mineral Physics I

## Chapter 4. Equation of state

### Section 8. Shock compression

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TOMOO KATSURA

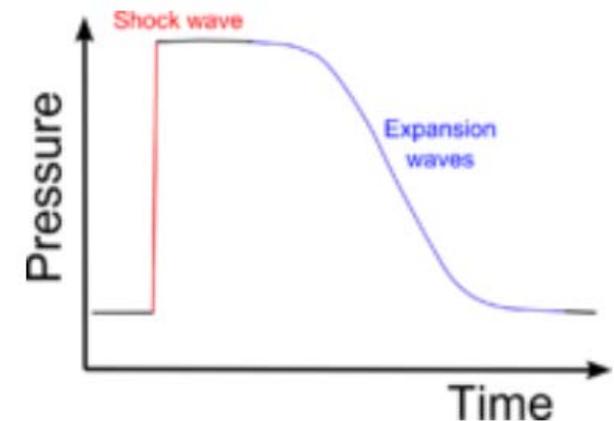
BAYERISCHES GEOINSTITUT, UNIVERSITY OF BAYREUTH,  
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# Shock wave

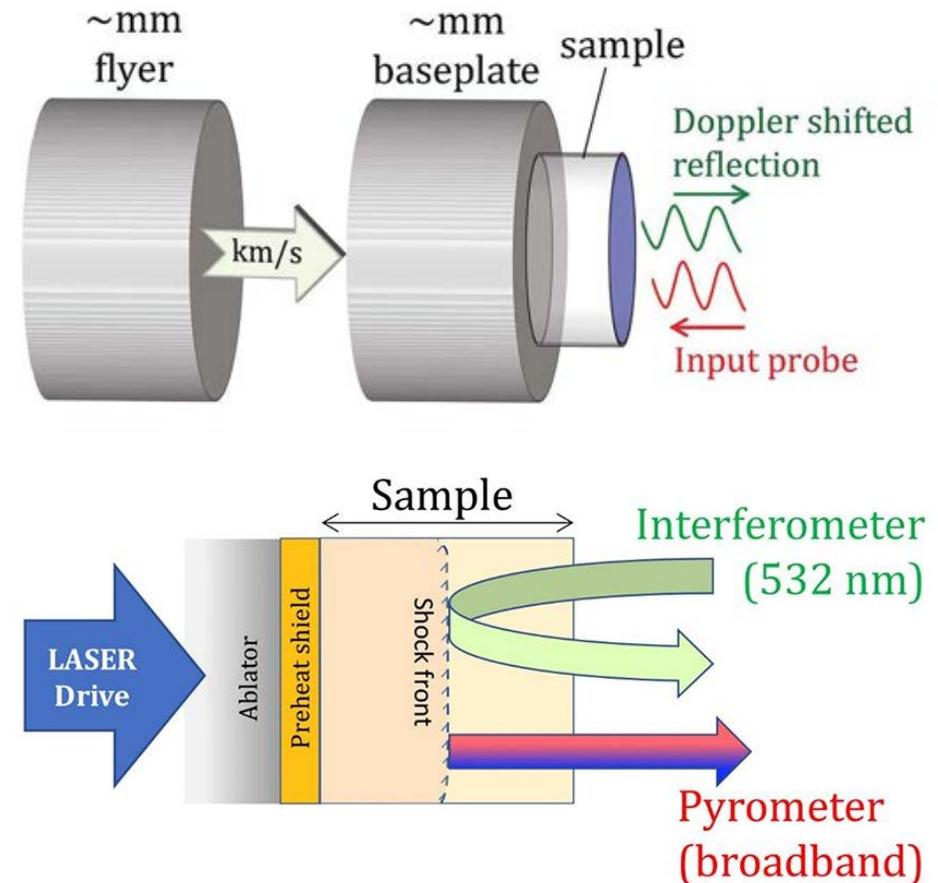
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- ❑ Shock compression: an experiment to make a body under extremely high pressure using shock wave
- ❑ **Shock wave:**
  - A type of propagating disturbance that moves **faster** than the **local speed of sound** in the medium.
    - ✓ ex. associated with the passage of the ultrasonic aircraft and missile
  - Characterized by an abrupt, nearly discontinuous change in pressure, temperature and density of the medium



# Shock compression

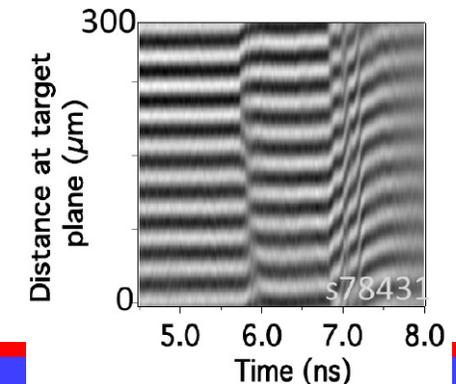
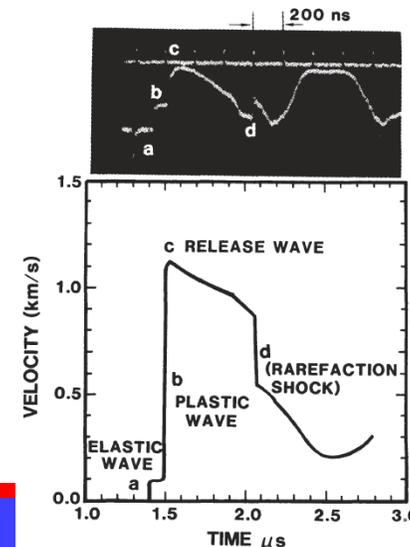
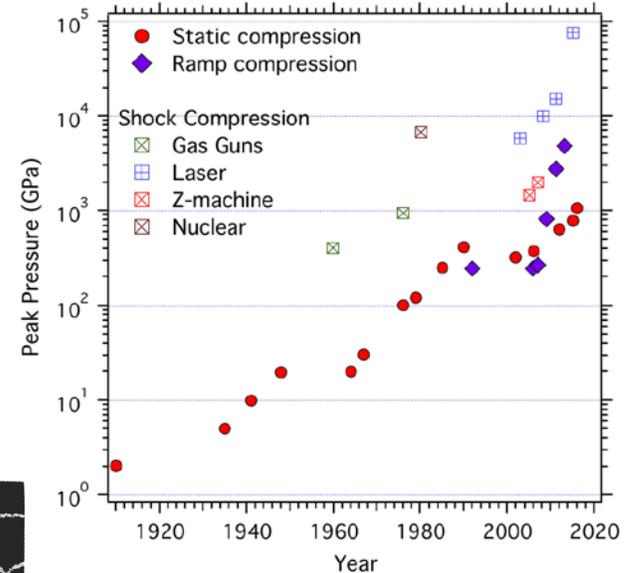
- ❑ Compress a solid or liquid by using a shock wave
  - **Explosion**
  - **Gas-gun**
    - ✓ Conventional technique
    - ✓ Collision of a **high-velocity flyer** accelerated by a gas-gun to a sample using a gun
  - **Laser ablation**
    - ✓ New technique
    - ✓ **Instantaneous heating** of the ablator by laser → **sudden evaporation** of the ablator → shock



# $P$ , $T$ , $t$ conditions of shock compression

- ❑  $P$  range: extremely high
  - Static compression using DAC
    - ✓ Can be 1 TPa, but mostly up to 200 GPa
  - Shock compression using a gas gun
    - ✓ Up to 1 TPa
  - Shock compression using the laser ablation
    - ✓ ~ 10 TPa or more
- ❑  $T$  range: can be high according to the pressure and material
  - Can be 10,000 K in the TPa range
  - A large proportion of shock energy is used to increase the temperature in addition to the pressure
- ❑ Duration of high pressure: depending on the pressure and the size of the sample
  - a few  $\mu\text{s}$  ~ 1 ns

Duffy & Smith (2019)



# Advantage and disadvantage of shock experiment

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## □ Advantages

- Can generate extremely pressures
- Can evaluate pressure without any pressure scale
- Can evaluate density change

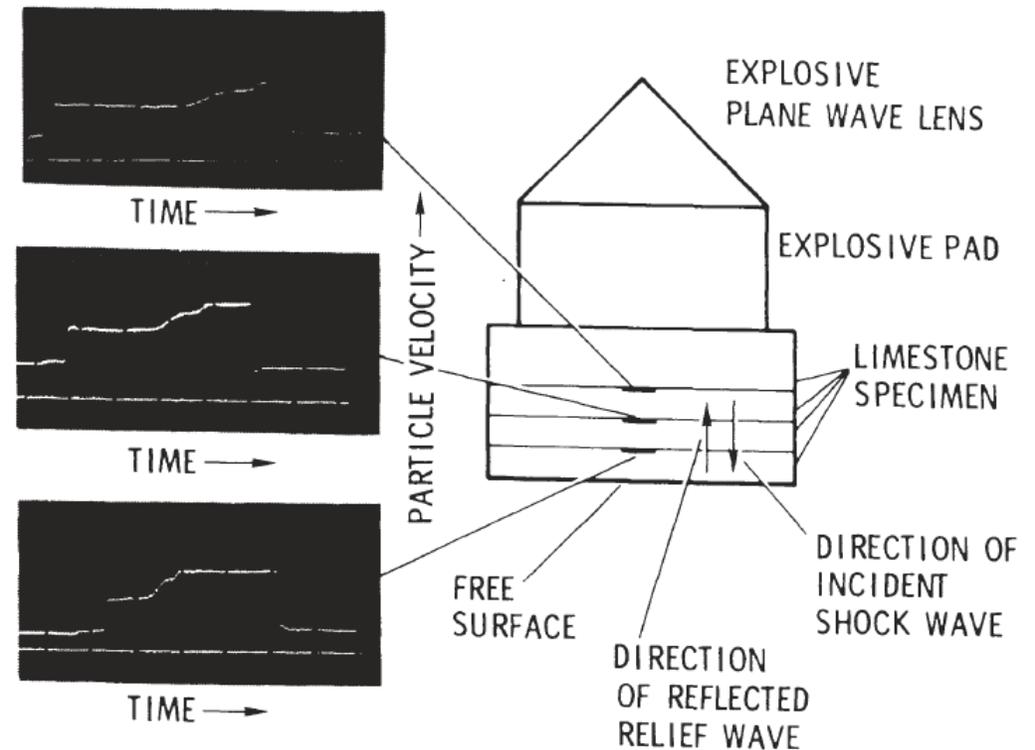
## □ Disadvantages

- Very short ( $\mu\text{s}$  to  $\text{ns}$ ) time scale
- Temperature cannot be controlled independently from the pressure
  - ✓ Spontaneous temperature increase due to results of shock compression
- Large errors in temperature estimation



# Gas-gun type shock compression

- ❑ This lecture: typically explain the gas-gun type shock compression experiment
  - Conventional, and therefore easy to understand
- ❑ Measured quantities
  - Velocity of the flyer (projectile)
    - ✓ “particle velocity”
    - ✓ Regarded as the velocity of motion of atoms or particles composing the material
  - Shock velocity



# An example of shock compression data

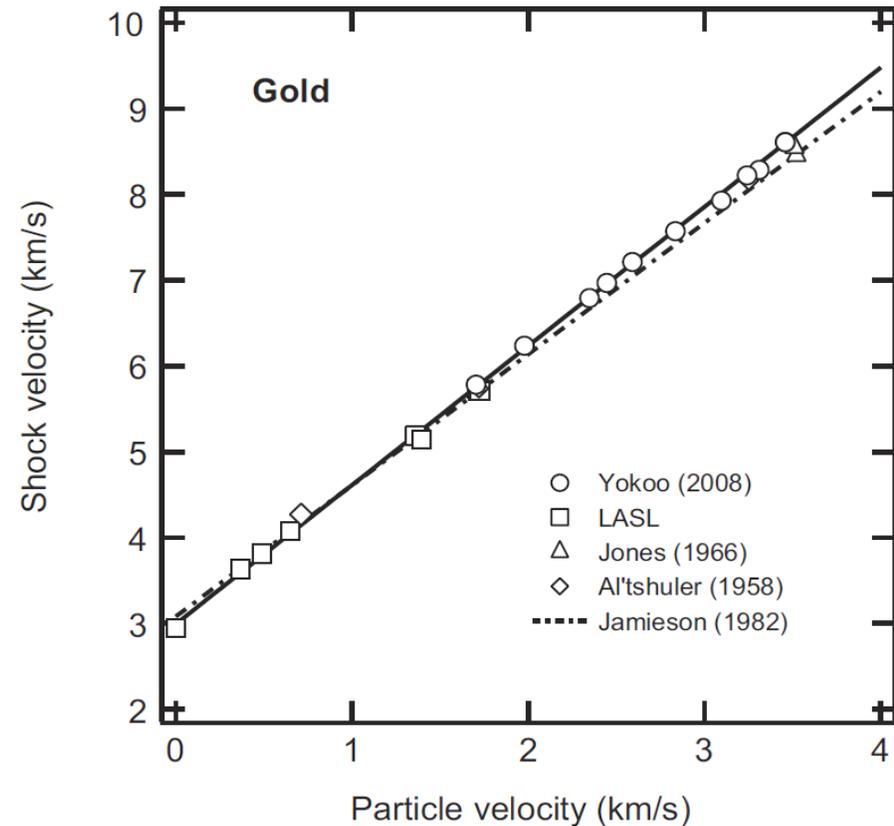
## □ Example: gold

➤  $U_s = 2.995 + 1.653u_p - 0.013u_p^2$  (4.8.1)

✓  $U_s$ : shock velocity

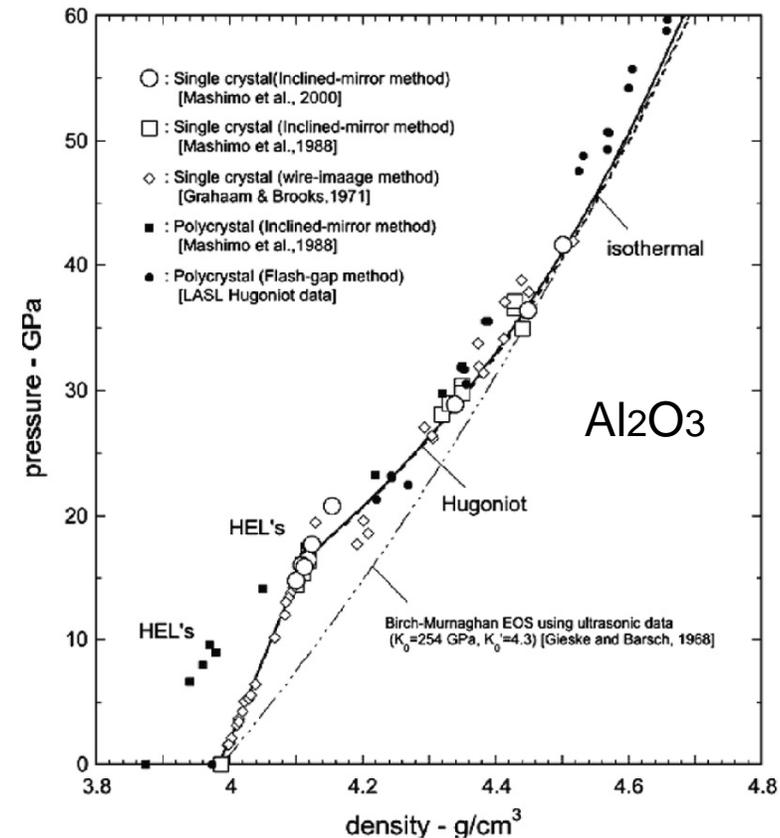
✓  $u_p$ : particle velocity

- Los Alamos Shock Hugoniot Data, edited by S. P. Marsh, University of California Press, Berkeley, 1979; A. H. Jones *et al.*, *J. Appl. Phys.* 37, 3493, 1966; L. V. Al'tshuler *et al.*, *J. Appl. Mech. Tech. Phys.* 22, 145, 1981; M. Yokoo *et al.*, *Appl. Phys. Lett.* 92, 051901, 2008.



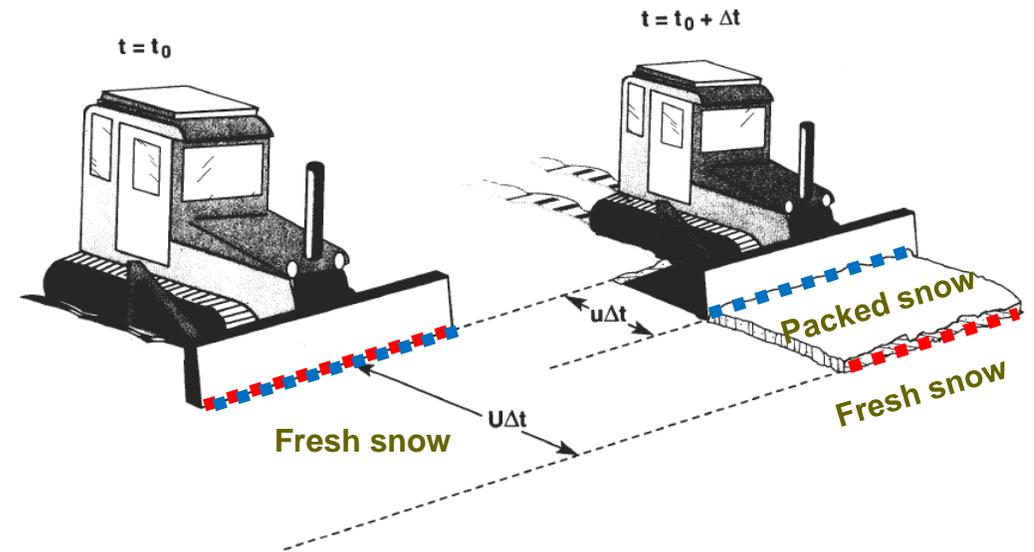
# Hugoniot elastic limit

- ❑ Solid: having significant to the shear
  - resist to shock wave up to some conditions
- ❑ Hugoniot elastic limit (HEL)
  - The matter behaves as a liquid above the HEL
  - The data below the HEL cannot be used to obtain an equation of state
  - Examples:
    - ✓  $\text{Al}_2\text{O}_3$ : about 20 GPa (very hard)
    - ✓ Au: almost none (very soft)



# Physical image of shock wave: analogy of a moving snowplow

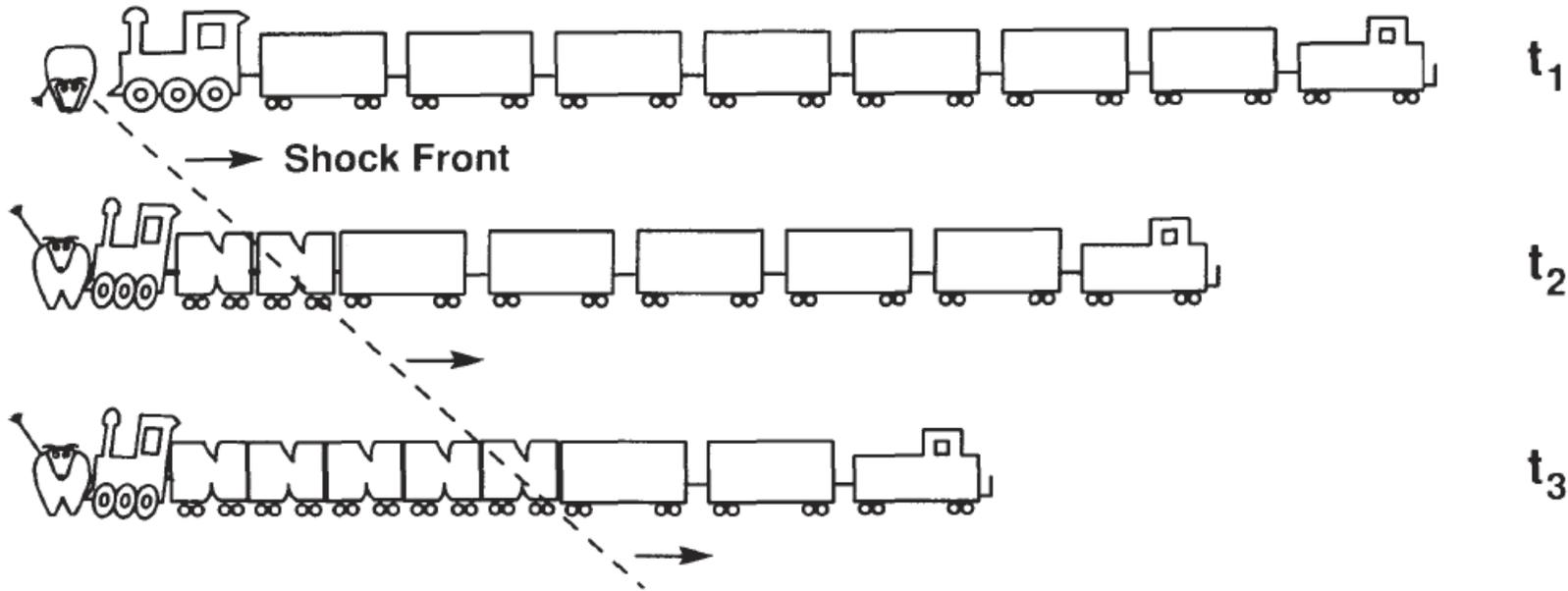
- ❑ Shock wave: Ubiquitous results of matters moving at velocities faster than the speeds at which adjacent material can move
- ❑ Example: the flow of snow in front of a moving snowplow
  - Movement of plow into fresh loose snow → formation of packed snow
  - ✓ The interface of fresh and packed snows moves faster than the plow



- Blue line: the front of plow
- Red line: the interface of the fresh and packed snows



# Physical image of shock wave: analogy of a moving train



Assuming that the bull is so strong that it stays the collision point.

- A bull is struck by a moving train
  - Shock wave: a moving discontinuity that separates undisturbed boxcars from crushed ones
  - Moving from left to right away from the point of collision



# Particle and shock velocities

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- ❑ Particle velocity: the velocity that a given element in the material acquires as a result of the shock wave passing over the element
  - We approximate that the elements have no velocity before the shock wave arrives.
  
- ❑ Shock velocity: the velocity with which the disturbance moves through the body
  
- ❑ Shock velocity  $\gg$  particle velocity
  - Ex. Debris such as dust or paper blown by a gust of wind
    - ✓ Shock velocity  $\rightarrow$  the wind velocity
    - ✓ Particle velocity  $\rightarrow$  the velocity of debris movement
    - ✓ The wind velocity  $\gg$  the debris velocity



# Equation for shock wave

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□ The fundamental points:

- Mass, momentum, and energy are conserved across the shock discontinuity
- → Three conservation equations
- → Rankine-Hugoniot equation

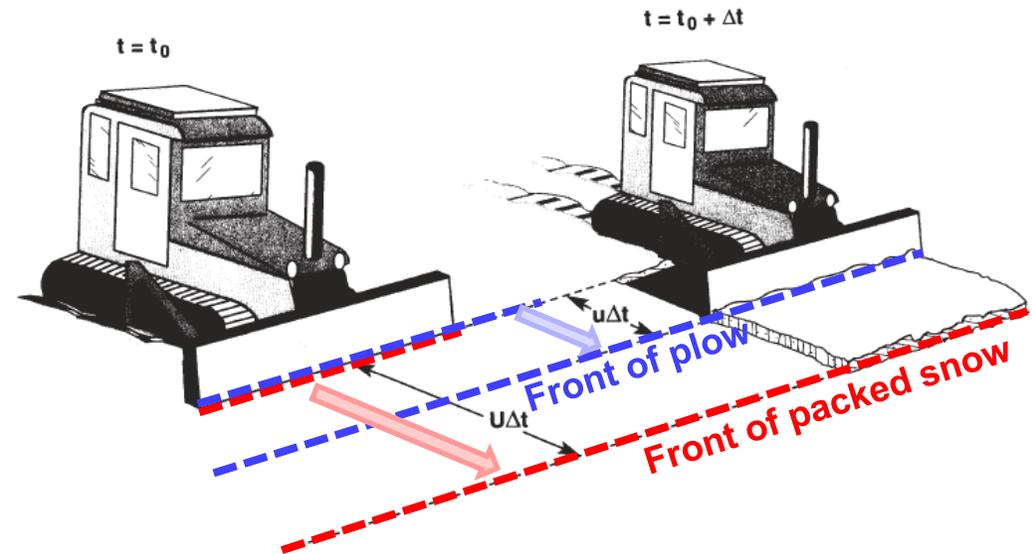
$$\checkmark E - E_0 = \frac{1}{2}(V_0 - V)(P + P_0) \quad (4.8.2)$$

- ✓ The relations among energy increase, volume change and shock pressure upon shock compression
- ✓ This equation will be derived from the next slide.



# Conservation of mass -1

- ❑ At time  $t_0$  the plow begins moving at velocity  $u$  from its starting point.
- ❑ At some later time,  $t_0 + \Delta t$ , the plow has moved a distance  $u\Delta t$ .
- ❑ The discontinuity between the loose and packed snow has moved a distance  $U\Delta t$ .
  - $U > u$ 
    - ✓ the discontinuity has advanced ahead of the plow



# Conservation of mass -2

1  
by



1  
the

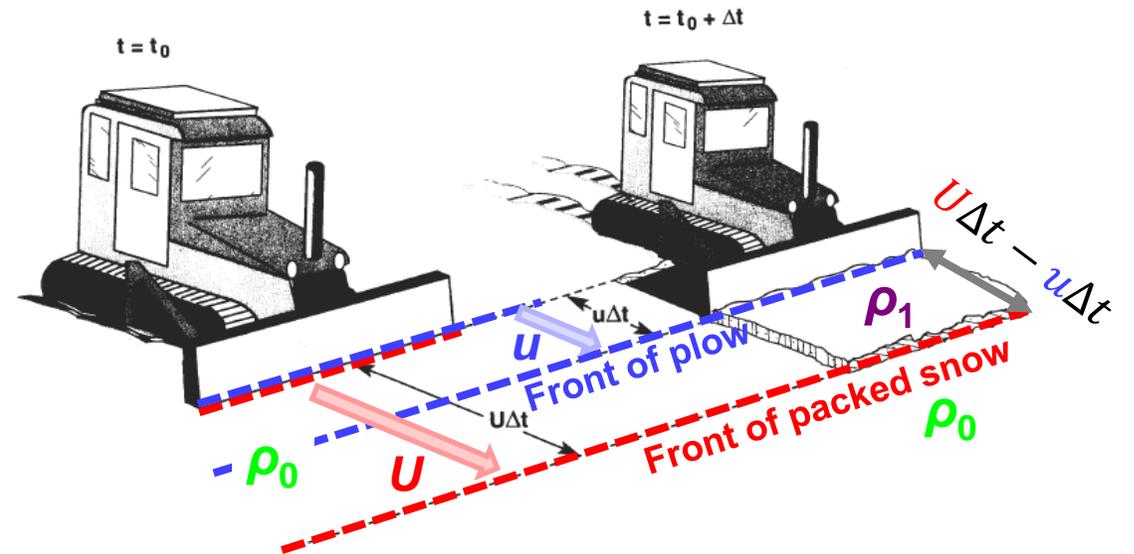


E



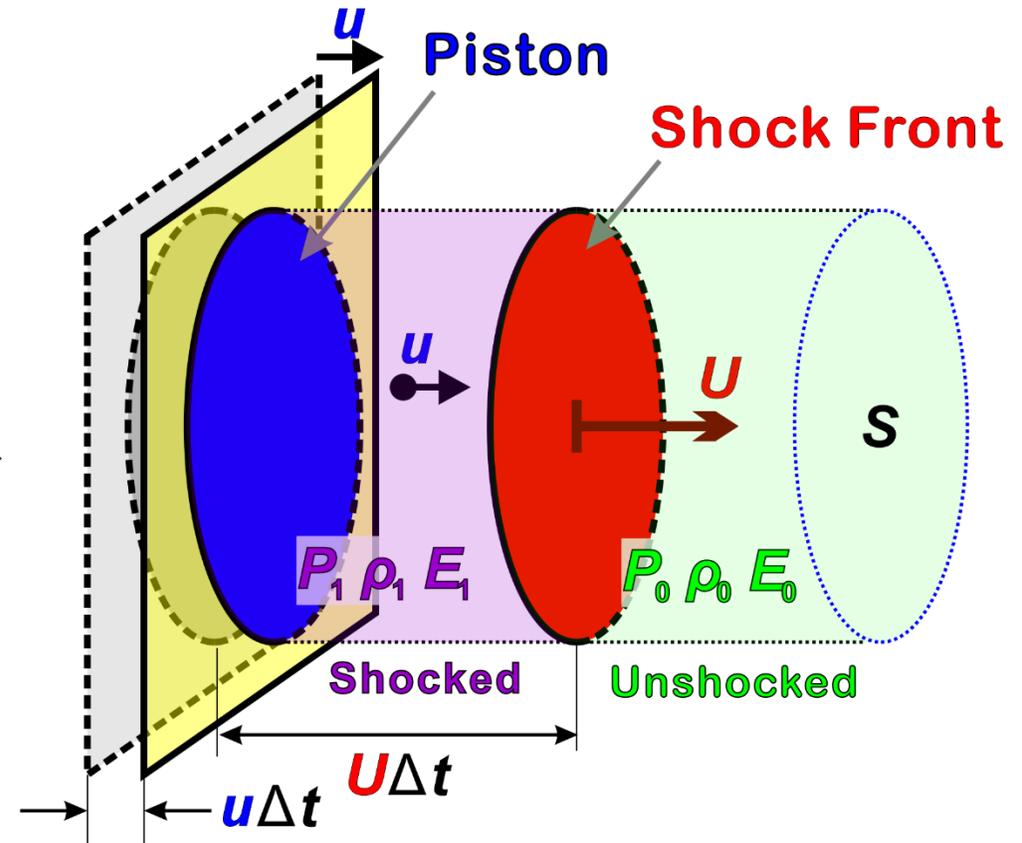
$\rho_0 U \Delta t = \rho_1 (U \Delta t - u \Delta t)$

✓ Conservation of mass



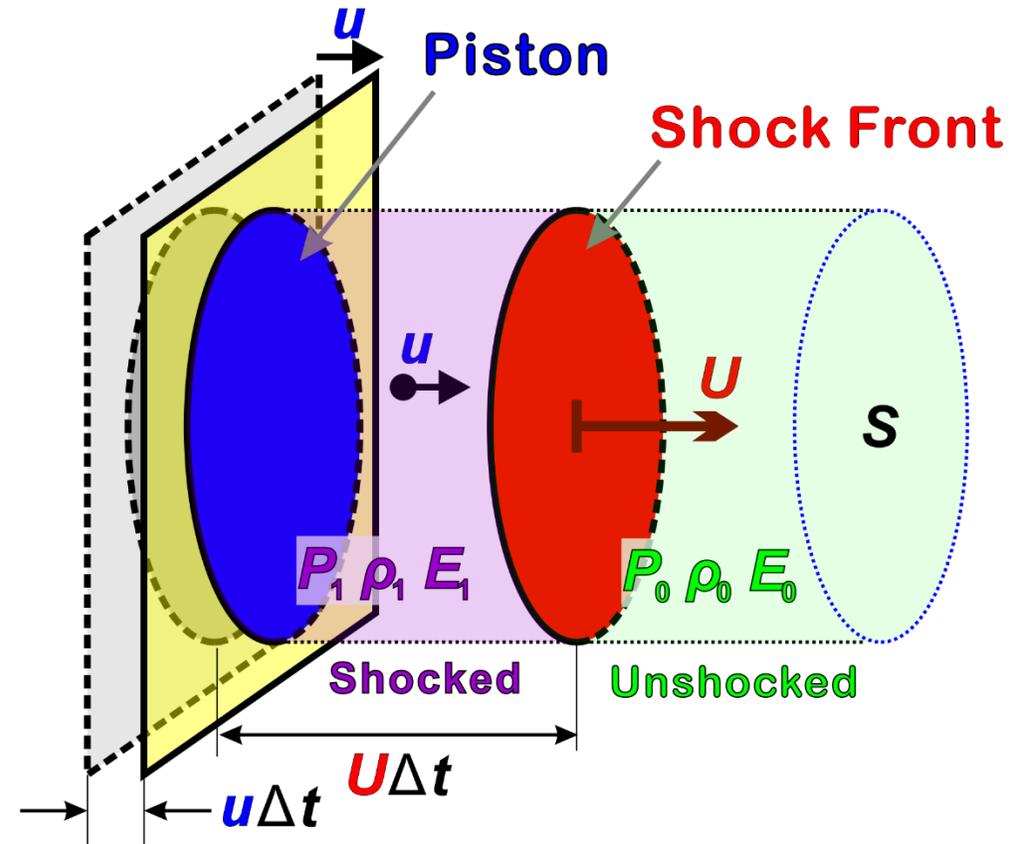
# Piston – fluid system

- ❑ Piston-fluid system
  - Replacing the plow-snow system
- ❑ Reference frame: undisturbed fluid
  - Zero velocity
- ❑ Movement of piston at velocity  $u$  → shock wave with velocity  $U$ 
  - Initial state “0” in front of the shock wave → state “1” behind the shock wave
- ❑ Particle velocity,  $u$ 
  - The velocity of a particle of the fluid caught up in the flow
  - The shocked material carried with it
  - piston velocity



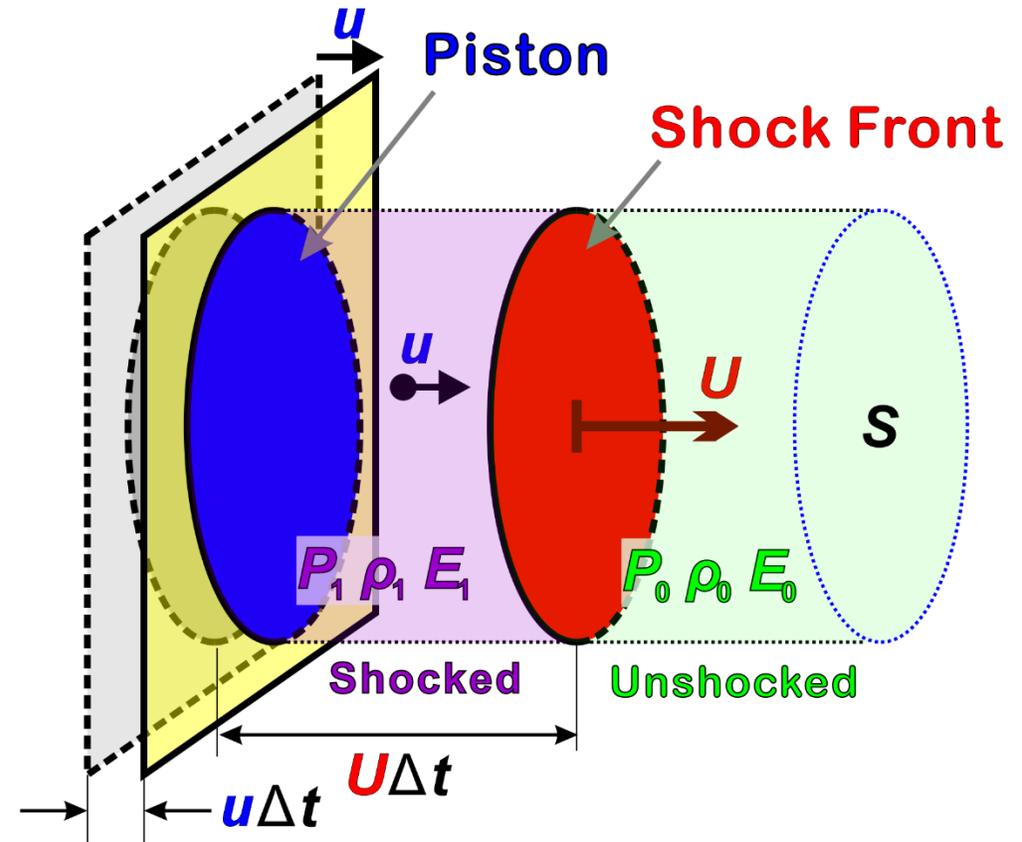
# Conservation of momentum -1

- Pressures before and after the shock wave
  - $P_0$ : the pressure in the unshocked part (before the shock)
    - ✓ from right to left
  - $P_1$ : the pressure applied by the piston (after the shock)
    - ✓ from left to right
  - The net force on the system:  
 $(P_1 - P_0)S$
  - The increase in momentum for  $\Delta t$ :  
 $(P_1 - P_0)S\Delta t$



# Conservation of momentum -2

- Increase in momentum for  $\Delta t$ 
  - The distance of the shock wave passing over:  $U\Delta t$
  - Volume of the shock passing over:  $U\Delta t \cdot S$
  - Mass of this volume:  $\rho_0 \cdot U\Delta t S$
  - Increased momentum:  $\rho_0 U\Delta t S \cdot u$
- Conservation of momentum
  - $(P_1 - P_0)S\Delta t = \rho_0 U\Delta t S u$
  - $P_1 - P_0 = \rho_0 U u$  (4.8.4)



# Conservation of energy -1

- The work done by the piston for  $\Delta t$ :

$$W = Fx = P_1 S \cdot u\Delta t = P_1 u S \Delta t$$

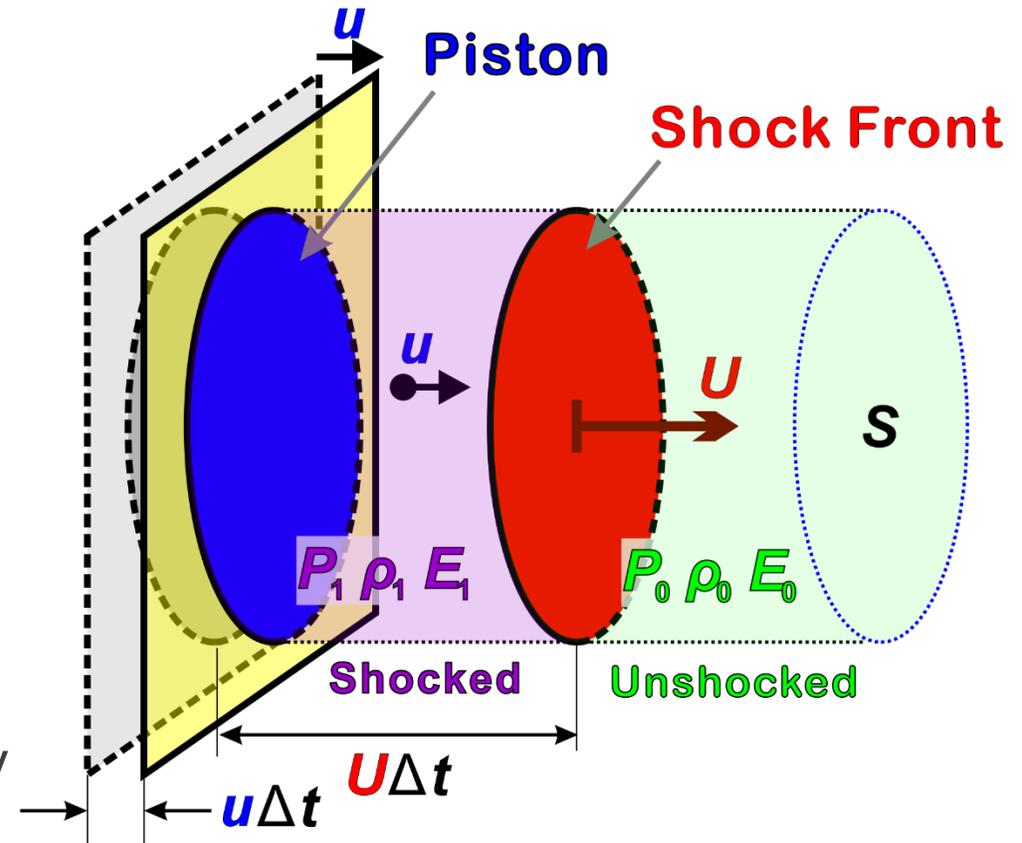
- Force:  $F = P_1 S$

- Piston movement:  $x = u\Delta t$

- The kinetic energy acquired by the volume where the shock wave has passed:

- $\frac{1}{2} m u^2 = \frac{1}{2} (\rho_0 U S \Delta t) u^2 = \frac{(\rho_0 U u^2) S \Delta t}{2}$   
(4.8.5)

- ✓ The mass of material accelerated by the shock wave:  $m = \rho_0 U S \Delta t$



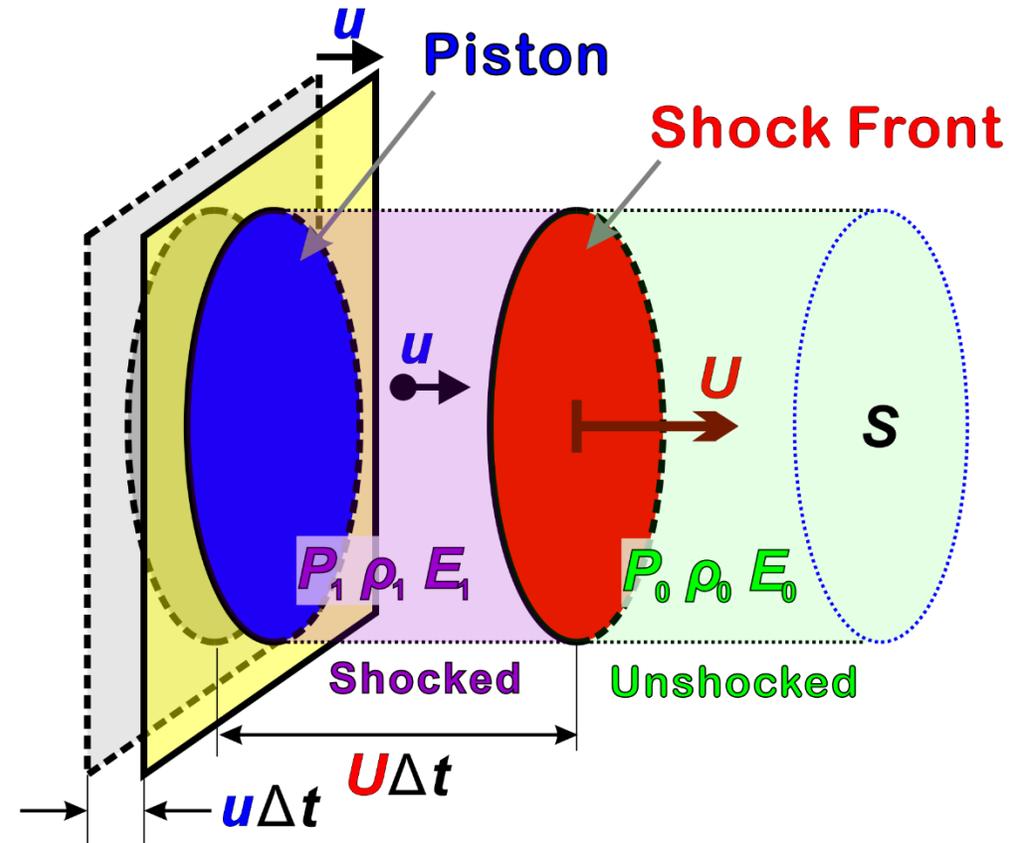
# Conservation of energy -2

□ The work done by the piston is also used to increase the internal energy in the volume of the shock wave passing over:  $(E_1 - E_0)\rho_0 U S \Delta t$

➤  $E_1, E_0$ : internal energy per mass

□ The work done by the shock  $P_1 u \Delta t$  becomes kinetic energy  $\frac{(\rho_0 U u^2) S \Delta t}{2}$  and internal energy  $(E_1 - E_0)\rho_0 U S \Delta t$

➤ 
$$P_1 u = \frac{1}{2} \rho_0 U u^2 + \rho_0 U (E_1 - E_0)$$
 (4.8.6)



# Rankine-Hugoniot equation -1

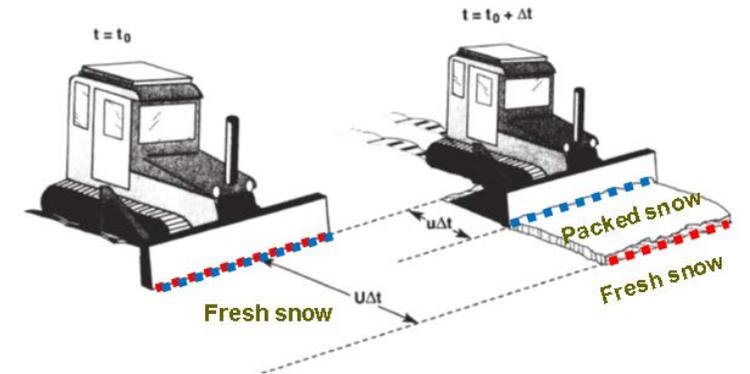
□ The  $E$  change with respect to  $P$  and specific  $V$  ( $V_1 \equiv 1/\rho_1$  and  $V_0 \equiv 1/\rho_0$ ) from the initial to final states by eliminating the  $U$  and  $u$  from Eq. (4.8.3), (4.8.4) and (4.8.6)

$$\triangleright (4.8.3): \rho_0 U = \rho_1 (U - u) \rightarrow \rho_1 U - \rho_0 U = \rho_1 u$$

$$\checkmark \rightarrow \left( \frac{1}{V_1} - \frac{1}{V_0} \right) U = \frac{1}{V_1} u \rightarrow \frac{V_0 - V_1}{V_0 V_1} U = \frac{1}{V_1} u$$

$$\checkmark \rightarrow U = \frac{V_0}{V_0 - V_1} u \quad (4.8.7)$$

- If the volume becomes zero after the shock wave passing over ( $V_1 = 0$ ),  $U = u$ 
  - No packed snow in front of the snowplow.
- If the volume does not decrease by the shock wave passing over ( $V_1 = V_0$ ),  $U = \infty$ 
  - The shock front propagates infinitely fast



If  $\rho_0 = \rho_1$ , the packed snow front is equal to the end of fresh snow.



# Rankine-Hugoniot equation -2

□ Conservation of momentum (4.8.4)  $P_1 - P_0 = \rho_0 U u$  ← Conservation of mass

(4.8.7)  $U = \frac{V_0}{V_0 - V_1} u$

➤  $P_1 - P_0 = \rho_0 \left( \frac{V_0}{V_0 - V_1} u \right) u = \frac{\rho_0 u^2}{V_0 - V_1}$

➤  $\rightarrow u^2 = (P_1 - P_0)(V_0 - V_1)$  (4.8.8)

□ Conservation of energy (4.8.6)  $P_1 u = \frac{1}{2} \rho_0 U u^2 + \rho_0 U (E_1 - E_0) U u$  ←

Conservation of mass (4.8.7)  $U = \frac{V_0}{V_0 - V_1} u$

➤  $P_1 u = \frac{1}{2} \left( \rho_0 \frac{V_0 u}{V_0 - V_1} u^2 \right) + \rho_0 \frac{V_0 u}{V_0 - V_1} (E_1 - E_0)$

➤  $\rightarrow P_1 (V_0 - V_1) = \frac{1}{2} \rho_0 u^2 + \rho_0 (E_1 - E_0) (V_0 - V_1)$  (4.8.9)



# Rankine-Hugoniot equation -3

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$$\square (4.8.9) \quad P_1(V_0 - V_1) = \frac{1}{2}u^2 + (E_1 - E_0) \leftarrow (4.8.8) \quad u^2 = (P_1 - P_0)(V_0 - V_1)$$

$$\triangleright P_1(V_0 - V_1) = \frac{1}{2}[u^2] + (E_1 - E_0) = \frac{1}{2}[(P_1 - P_0)(V_0 - V_1)] + (E_1 - E_0)$$

$$\triangleright E_1 - E_0 = P_1(V_0 - V_1) - \frac{1}{2}(P_1 - P_0)(V_0 - V_1)$$

$$\triangleright E_1 - E_0 = \frac{1}{2}(V_0 - V_1)(P_1 + P_0) \quad (4.8.2)$$

## ✓ Rankine-Hugoniot equation

- Jump conditions
- The **initial** and **final** states achieved through the shock transition must be states of mechanical equilibrium



# $P, V, E$ at the shock compression

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## □ Final volume

➤ Conservation of mass (4.8.3)  $\rho_0 U = \rho_1 (U - u)$

➤  $\rightarrow \frac{\rho_0}{\rho_1} = \frac{V_1}{V_0} = \frac{U-u}{U} \rightarrow V_1 = V_0 \frac{U-u}{U}$  (4.8.3')

## □ Final pressure

➤ Conservation of momentum (4.8.4)  $P_1 - P_0 = \rho_0 U u$

➤  $\rightarrow P_1 = P_0 + \rho_0 U u = P_0 + \frac{U u}{V_0}$  (4.8.4')

## □ Final internal energy

➤ Rankine-Hugoniot equation (4.8.2)  $E_1 - E_0 = \frac{1}{2} (V_0 - V_1) (P_1 + P_0)$

➤  $E_1 = E_0 + \frac{1}{2} \left( V_0 - V_0 \frac{U-u}{U} \right) \left( P_0 + \frac{U u}{V_0} + P_0 \right) = E_0 + V_0 \left( 1 - \frac{U-u}{U} \right) \left( P_0 + \frac{1}{2} \frac{U u}{V_0} \right)$  (4.8.2')



# $T$ at the shock compression

$$\square E_1 - E_0 = \Delta E_{T_0} + \Delta E_{\text{th}} \quad (4.8.10)$$

✓  $\Delta E_{T_0}$ : Internal energy increase by compression from  $V_0$  to  $V_1$  at  $T_0$

✓  $\Delta E_{\text{th}}$ : internal energy increase by heating from  $T_0$  to  $T_1$  at  $V_1$

$$\square \Delta E_{\text{th}} = \int_{T_0}^{T_1} C_V(V_1) dT = 9nR \left( \frac{T}{\Theta_D(V_1)} \right)^3 \int_{\frac{\Theta_D(V_1)}{T_0}}^{\frac{\Theta_D(V_1)}{T_1}} \frac{x^4 e^x}{(e^x - 1)^2} dT \quad (4.8.11)$$

$$\triangleright \Theta_D(V_1) = \Theta_D(V_0) \exp \left[ \frac{\gamma_0}{q} \left\{ 1 - \left( \frac{V_1}{V_0} \right)^q \right\} \right] \quad (4.8.12)$$

$$\square \text{The 1}^{\text{st}} \text{ and 2}^{\text{nd}} \text{ laws of thermodynamics (1.1.5) } dE = TdS - PdV$$

$$\triangleright \Delta E_{T_0} = \int_{V_0}^{V_1} [T_0 dS - P_{T_0}(V) dV] = T_0 [S(T_0, V_1) - S(T_0, V_0)] - \int_{V_0}^{V_1} P_{T_0}(V) dV \quad (4.8.13)$$

$$\checkmark S(T_0, V_1) = \int_0^{T_0} \frac{C_V(V_1)}{T} dT, \quad S(T_0, V_0) = \int_0^{T_0} \frac{C_V(V_0)}{T} dT \quad (4.8.14)$$

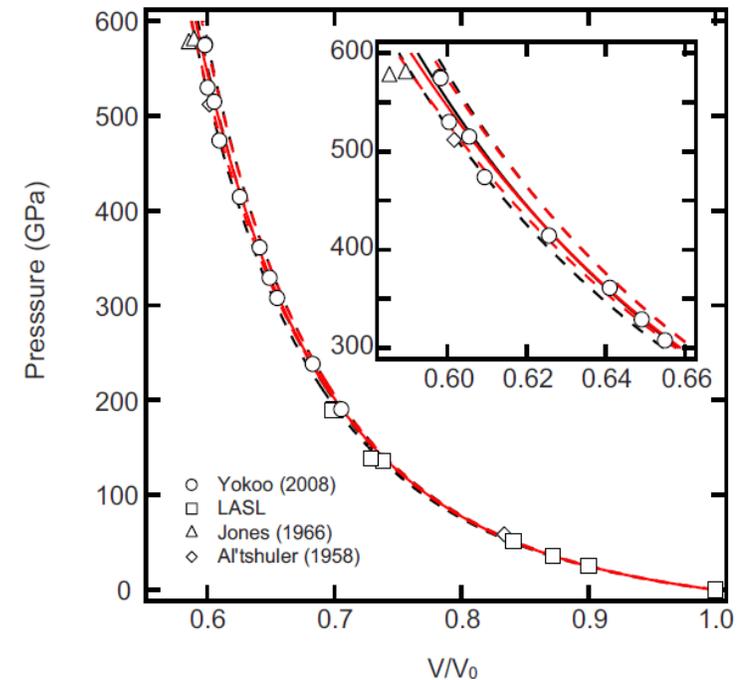
✓  $\int_{V_0}^{V_1} P_{T_0}(V) dV$ : along the isothermal equation of state

$$\square T_1 \text{ is inversely calculated from Eq. (4.8.10) to (4.8.14)}$$



# Hugoniot

- Three equations (4.8.2, 3, 6) and one static EOS for five variables ( $P$ ,  $V$ ,  $E$ ,  $U$  and  $u$ )
  - Only one independent variable
  - One curve in a  $P$ - $V$ - $E$ - $U$ - $u$  space
    - ✓ “**Rankine-Hugoniot curve**” or simply “**Hugoniot**”
    - ✓ Can be represented in any 2D plane ( $P$ - $V$ ,  $P$ - $u$ ,  $U$ - $u$ ,...)
    - ✓ Not a path that is followed during compression or any special thermodynamic path
    - ✓ The locus of all the possible end states that can be achieved behind a single shock wave passing through a material at a given initial state.



Gold (Yokoo et al., 2009)



# Rayleigh Line

□ Rayleigh Line ( $\mathcal{R}$ ): a line connecting the initial and final shocked state

□ From (4.8.8)  $u^2 = (P_1 - P_0)(V_0 - V_1)$  with (4.8.7)  $U = \frac{V_0}{V_0 - V_1} u$

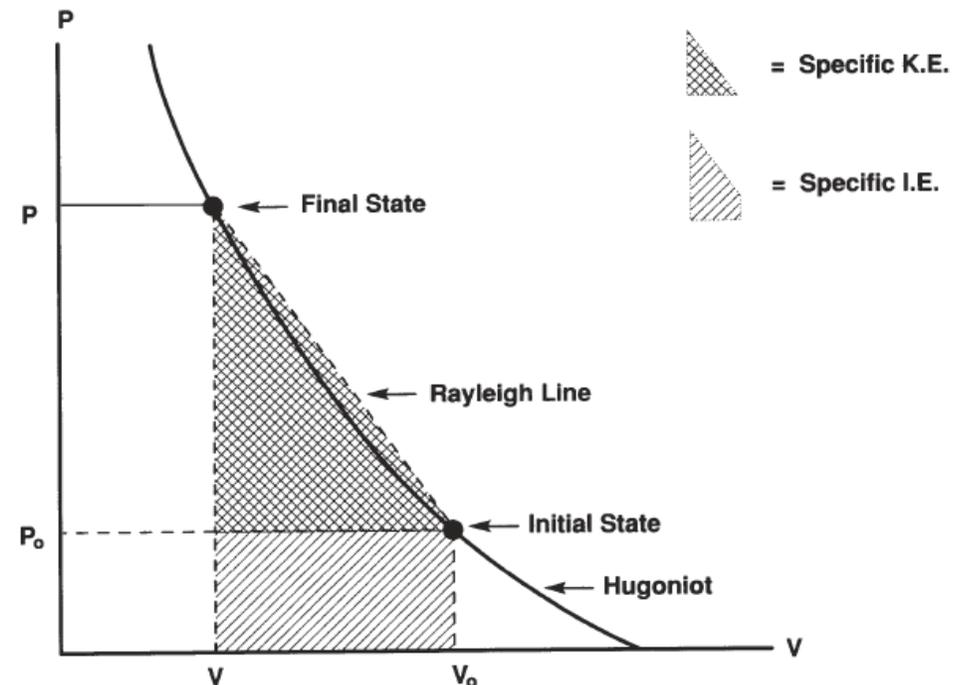
$$\text{➤ } u = \sqrt{(P_1 - P_0)(V_0 - V_1)}$$

$$\text{➤ } U = \frac{V_0}{V_0 - V_1} u = \sqrt{\frac{P_1 - P_0}{V_0 - V_1}} \quad (4.8.15)$$

□ The slope of  $\mathcal{R} = \frac{P_1 - P_0}{V_1 - V_0} = -\left(\frac{U}{V_0}\right)^2 \quad (4.8.16)$

➤ With increasing  $U$ , the slope becomes steeper rapidly.

✓ Hugoniot: concave downward



# Kinetic and internal energy increase by shock compression

□ Kinetic energy from (4.8.8)

$$\triangleright \frac{1}{2}u^2 = \frac{1}{2}(P - P_0)(V_0 - V)$$

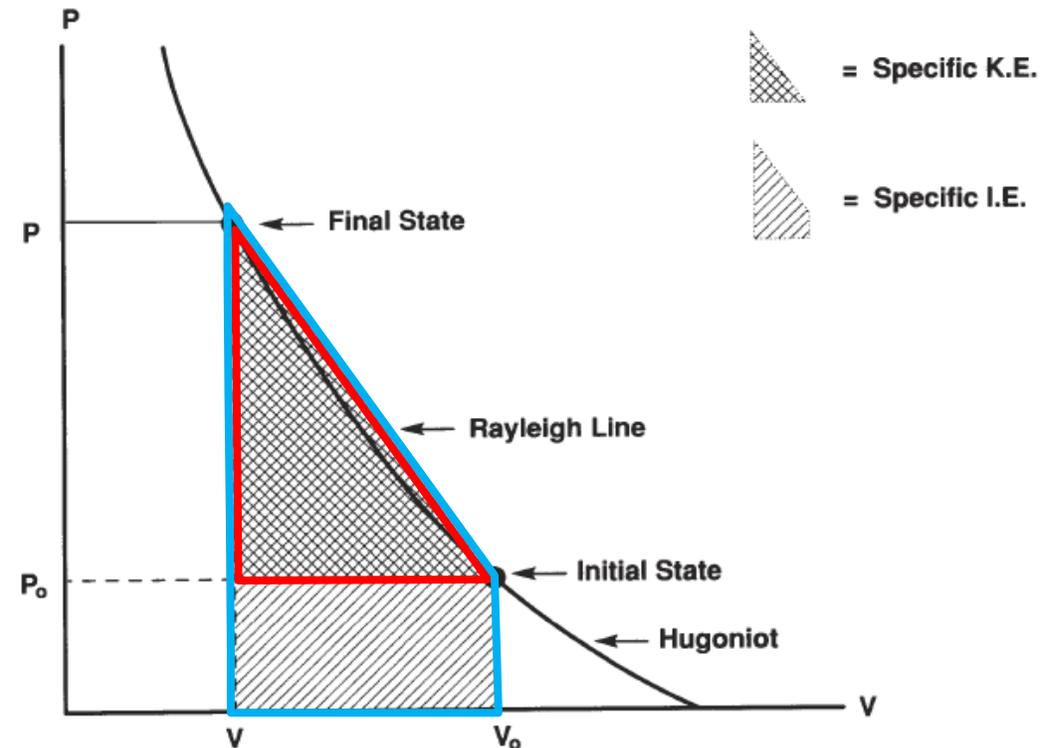
□ Internal energy increase from the Rankine-Hugoniot equation (4.8.2)

$$\triangleright E - E_1 = \frac{1}{2}(P + P_0)(V_0 - V)$$

□ If initially zero pressure ( $P_0 = 0$ ),

$$\triangleright E - E_1 = \frac{1}{2}P(V_0 - V) = \frac{1}{2}u^2 \quad (4.8.11)$$

➤ The internal energy increase is equal to the kinetic energy increase



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End

